

ON UNCONDITIONAL WELL-POSEDNESS OF MODIFIED KDV

SOONSIK KWON AND TADAHIRO OH

ABSTRACT. Bourgain [2] proved that the periodic modified KdV equation (mKdV) is locally well-posed in $H^s(\mathbb{T})$, $s \geq 1/2$, by introducing new weighted Sobolev spaces $X^{s,b}$, where the uniqueness holds *conditionally*, namely in $C([0, T]; H^s) \cap X^{s, \frac{1}{2}}([0, T] \times \mathbb{T})$. In this paper, we establish *unconditional* well-posedness of mKdV in $H^s(\mathbb{T})$, $s \geq \frac{1}{2}$, i.e. in addition we establish *unconditional* uniqueness in $C([0, T]; H^s)$, $s \geq 1/2$, of solutions to mKdV. We prove this result via differentiation by parts. For the endpoint case $s = \frac{1}{2}$, we perform careful quinti- and septi-linear estimates after the second differentiation by parts.

CONTENTS

1. Introduction	1
2. Nonlinear estimates for $s > \frac{1}{2}$	6
3. Unconditional local well-posedness for $s > \frac{1}{2}$	9
4. Endpoint case: $s = \frac{1}{2}$	11
References	17

1. INTRODUCTION

We consider the Cauchy problem for the modified Korteweg-de Vries (mKdV) equation on the one-dimensional torus $\mathbb{T} = \mathbb{R}/2\pi\mathbb{Z}$:

$$\begin{cases} \partial_t u = \partial_x^3 u \pm u^2 \partial_x u, \\ u|_{t=0} = u_0, \end{cases} \quad (x, t) \in \mathbb{T} \times \mathbb{R} \quad (1.1)$$

where u is a real-valued function. The mKdV has received a great deal of attention both from applied and theoretical fields and is known to be completely integrable in the sense that it enjoys the Lax pair structure and so infinitely many conservation laws. In particular, if u is a “nice” solution of (1.1), then the L^2 -norm is conserved. i.e. $\|u(t)\|_{L^2} = \|u_0\|_{L^2}$. Then, by the change of variables $x \rightarrow x \mp \mu t$ with $\mu = \frac{1}{2\pi} \|u_0\|_{L^2}^2$, we can rewrite (1.1) as

$$\begin{cases} \partial_t u = \partial_x^3 u \pm \left(u^2 - \frac{1}{2\pi} \int_{\mathbb{T}} u^2 \right) \partial_x u, \\ u|_{t=0} = u_0. \end{cases} \quad (1.2)$$

In this paper, we study the unconditional local well-posedness of (1.2).

Date: July 25, 2011.

2000 *Mathematics Subject Classification.* 35Q53.

Key words and phrases. modified KdV; well-posedness; uniqueness.

The first author is supported by NRF grant 2010-0024017.

Let us briefly go over recent results on the well-posedness theory of the periodic mKdV. In [2], Bourgain introduced a new weighted space-time Sobolev space $X^{s,b}$ (also known as *dispersive-Sobolev space*), whose norm is given by

$$\|u\|_{X^{s,b}(\mathbb{T} \times \mathbb{R})} = \|\langle k \rangle^s \langle \tau + k^3 \rangle^b \widehat{u}(k, \tau)\|_{l_k^2 L_\tau^2(\mathbb{Z} \times \mathbb{R})}, \quad (1.3)$$

where $\langle \cdot \rangle = 1 + |\cdot|$. By the fixed point argument in an appropriate $X^{s,b}$ space, he proved that (1.2) is locally well-posed in $H^s(\mathbb{T})$, $s \geq \frac{1}{2}$. Then, Colliander-Keel-Staffilani-Takaoka-Tao [6] proved global well-posedness in H^s , $s \geq \frac{1}{2}$, via the I -method. We point out that the solution map $\mathcal{S}_t : u_0 \in H^s \mapsto u(t) \in H^s$ constructed in [2, 6] is smooth. Indeed, it was shown in [3] that the solution map to (1.2) can not be smooth in H^s for $s < \frac{1}{2}$. See [5, 12] for related results. Nonetheless, Takaoka-Tsutsumi [12] successfully modified the $X^{s,b}$ space to reduce the nonlinear effect from the resonant term (see \mathcal{R} in (1.7) below) and proved local well-posedness in H^s , $s > \frac{3}{8}$. Nakanishi-Takaoka-Tsutsumi [11] further improved the result and proved local well-posedness in $H^s \cup \mathcal{FL}^{\frac{1}{2}, \infty}$, $s > \frac{1}{3}$, where $\|f\|_{\mathcal{FL}^{\frac{1}{2}, \infty}} = \sup_k \langle k \rangle^{\frac{1}{2}} |\widehat{f}(k)| < \infty$. (Existence alone holds in H^s for $s > \frac{1}{4}$.) Note that the solution map constructed in [12, 11] is not uniformly continuous. There is also a result using the complete integrability of the equation. Kappeler-Topalov [9] proved that defocusing mKdV, (1.1) with the $-$ sign, is globally well-posed in $L^2(\mathbb{T})$ via the inverse spectral method.

Now, let us examine the uniqueness of solutions in the above results. In [2, 12], the uniqueness (with prescribed L^2 -norm) holds in (a ball in) $C([0, T]; H^s) \cap X$, where X is an auxiliary function space, i.e. only within the (modified) $X^{s,b}$ space. Thus, the uniqueness holds *conditionally*, since uniqueness may not hold without the restriction of the auxiliary space X . In [9], the uniqueness holds in the class of solutions obtained by a limiting procedure of smooth solutions.

Recall the following definition from Kato [10]. We say that a Cauchy problem is *unconditionally well-posed* in H^s if for every initial condition $u_0 \in H^s$, there exist $T > 0$ and a *unique* solution $u \in C([0, T]; H^s)$ such that $u(0) = u_0$. Also, see [7]. We refer to such uniqueness in $C([0, T]; H^s)$ without intersecting with any auxiliary function space as *unconditional uniqueness*. Unconditional uniqueness is a concept of uniqueness which does not depend on how solutions are constructed. See, for example, Zhou [14] for unconditional uniqueness of KdV in $L^2(\mathbb{R})^1$ or Tao [13] for focusing mass-critical NLS with spherical symmetry in $L^2(\mathbb{R}^d)$, $d \geq 5$.

The main result of the paper is the following:

Theorem 1.1. *Let $s \geq \frac{1}{2}$. Then, mKdV is unconditionally locally well-posed in $H^s(\mathbb{T})$.*

Our result provides another proof of the local well-posedness. We think that this proof is more natural and elementary since we do not use any auxiliary function spaces but only rely on simple *differentiation by parts* and *Cauchy-Schwarz inequality*. As a result, we can establish unconditional uniqueness of solution to mKdV in $H^s(\mathbb{T})$, $s \geq \frac{1}{2}$, which is an improvement of Bourgain's result [2] in the aspect of uniqueness.

Remark 1.2. In the proof of Theorem 1.1, we show existence and uniqueness of solutions to the renormalized mKdV (1.2) in $H^s(\mathbb{T})$, $s \geq \frac{1}{2}$. From this, one can deduce existence and uniqueness of solutions to the original mKdV (1.1) in $H^s(\mathbb{T})$, $s \geq \frac{1}{2}$. Indeed, for a given $u_0 \in H^s(\mathbb{T})$ with $s \geq \frac{1}{2}$, a function $u(x, t) \in L^\infty([0, T]; H^s)$ is a solution to the original

¹This is uniqueness of weak solutions in the space of $L_t^\infty L_x^2(\mathbb{R} \times \mathbb{R})$

mKdV (1.1) with the initial condition u_0 if and only if v is a solution to the renormalized mKdV (1.2) with the same initial condition u_0 , where v is defined by

$$v(x, t) := u\left(x \mp \frac{1}{2\pi} \int_0^t \|u(t')\|_{L^2}^2 dt', t\right). \quad (1.4)$$

(Here, we used the fact that $\|v(t)\|_{L^2} = \|u(t)\|_{L^2}$ for v defined in (1.4).)

Now, suppose that u_1 and u_2 are two solutions to the original mKdV (1.1) in $C([0, T]; H^s)$ with the same initial condition $u_0 \in H^s(\mathbb{T})$ with $s \geq \frac{1}{2}$. Then, $t \mapsto \|u_j(t)\|_{L^2}^2$ is locally integrable and v_1 and v_2 defined via (1.4) are solutions to the renormalized mKdV (1.2) in $C([0, T]; H^s)$ with the same initial condition u_0 . Hence, by Theorem 1.1, we have $v_1 = v_2$ in $C([0, T]; H^s)$. In particular, $\|v_j(t)\|_{L^2}^2$, $j = 1, 2$, is constant in time.² In view of (1.4), we see that $\|u_j(t)\|_{L^2}^2$, $j = 1, 2$, is also constant in time, and the transformation (1.4) can be written as

$$v_j(x, t) = u_j\left(x \mp \frac{t}{2\pi} \|u_0\|_{L^2}^2, t\right), \quad \text{for } j = 1, 2. \quad (1.5)$$

Therefore, from (the inverse of) (1.5) and $v_1 = v_2$ in $C([0, T]; H^s)$, we obtain $u_1 = u_2$ in $C([0, T]; H^s)$. This shows unconditional uniqueness of the original mKdV (1.1).

Lastly, we discuss the regularity of the solution map: $u_0 \mapsto u(t)$ of the original mKdV (1.1) (for sufficiently small t depending on the size of initial data.) From the proof of Theorem 1.1, it follows that the solution map of the renormalized mKdV (1.2) is locally Lipschitz continuous. Consequently, this yields local Lipschitz continuity of the solution map of the original mKdV (1.1) in the class

$$\{u_0 \in H^s(\mathbb{T}) : \|u_0\|_{L^2(\mathbb{T})} = c\} \quad (1.6)$$

with a fixed c . (Two initial data of distinct L^2 -norms give rise to two different renormalized mKdV (1.2), and thus their solutions are not comparable. In general, one can show that the uniform continuity of the solution map of the original mKdV (1.1) fails without prescribing the L^2 -norm.)

Remark 1.3. Many of the unconditional uniqueness results use some auxiliary function spaces (e.g. $X^{s,b}$ spaces in [14], Strichartz spaces in [13]), which are designed to be large enough to contain $C([0, T]; H^s)$ such that desired nonlinear estimates hold. However, we simply use the $C([0, T]; H^s)$ -norm in the proof of Theorem 1.1.

We also point out that Theorem 1.1 does *not* imply unconditional well-posedness for KdV in $H^s(\mathbb{T})$, $s \geq -\frac{1}{2}$ even under the Miura map. Indeed, the issue of unconditional uniqueness of the periodic KdV is settled in view of the non-uniqueness result by Christ [4] for KdV in $C_t H^s(\mathbb{T})$, $s < 0$, and an (implicitly implied) positive result in $L^2(\mathbb{T})$ by Babin-Ilyin-Titi [1].

Theorem 1.1 with global well-posedness of mKdV in H^s , $s \geq \frac{1}{2}$, by [6] yields the following corollary.

Corollary 1.4. *Let $s \geq \frac{1}{2}$. Then, mKdV is unconditionally globally well-posed in $H^s(\mathbb{T})$.*

We prove Theorem 1.1 by establishing *a priori* estimates, where we use only the $C_t H_x^s$ -norm of solutions. In the following, we briefly describe the idea of differentiation by parts introduced in Babin-Ilyin-Titi [1].

Let $S(t) = e^{t\partial_x^3}$ denote the semigroup to the Airy equation (= linear part of mKdV (1.1).) We apply a change of coordinates: $v(t) = S(-t)u(t)$. In terms of the spatial

²For (unique) solutions to (1.2) in $C([0, T]; H^s)$, $s \geq \frac{1}{2}$, constructed in this paper, one can easily check that the L^2 -norm is conserved in time.

Fourier coefficients, this can be written as $v_k(t) = e^{ik^3 t} u_k(t)$, where $v_k(t)$ denotes the k -th (spatial) Fourier coefficient of $v(\cdot, t)$. i.e. $v_k(t) = \widehat{v}(k, t)$. Working in terms of v has certain advantages. Ginibre [8] says “In the language of Quantum Mechanics, this consists in working in the so-called interaction representation.” In [2], Bourgain made an effective use of this coordinate by introducing the $X^{s,b}$ spaces. From the definition (1.3), we have $\|u\|_{X^{s,b}} = \|v\|_{H_t^b H_x^s}$, i.e. a function u is in $X^{s,b}$ if and only if its interaction representation $v(t) = S(-t)u(t)$ is in the classical Sobolev space $H_t^b H_x^s$. A similar idea has been applied to study equations in hydrodynamics. See [1, Section 2] for a nice discussion.

With $v(t) = S(-t)u(t)$, it follows from (1.2) (see [2]) that v satisfies³

$$\begin{aligned} \partial_t v_k &= \sum_{\substack{k_1+k_2+k_3=k \\ k_1+k_2 \neq 0}} ik_3 e^{it\Phi(\bar{k})} v_{k_1} v_{k_2} v_{k_3} \\ &= \frac{i}{3} \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0}} ke^{it\Phi(\bar{k})} v_{k_1} v_{k_2} v_{k_3} - ik|v_k|^2 v_k =: \mathcal{N} + \mathcal{R}. \end{aligned} \quad (1.7)$$

where $\Phi(\bar{k}) = \Phi(k, k_1, k_2, k_3) := k^3 - k_1^3 - k_2^3 - k_3^3$. With $k = k_1 + k_2 + k_3$, we have

$$\Phi(\bar{k}) = 3(k_1 + k_2)(k_2 + k_3)(k_3 + k_1). \quad (1.8)$$

In this framework, the usual Duhamel formulation of (1.2) corresponds to

$$v(t) = v_0 + \int_0^t \mathcal{N}(v)(t') + \mathcal{R}(v)(t') dt', \quad (1.9)$$

where $\mathcal{N}(t')$ stands for $\mathcal{N}(v)(t') = \mathcal{N}(v(t'), v(t'), v(t'))$, and so on. Due to the presence of ∂_x in the nonlinearity, a direct estimate in H^s on the nonlinear part in (1.9) does not work.

Assume that v is smooth in the following, since our goal is to obtain a priori estimates on solutions. As in [1], we can differentiate \mathcal{N} by parts, Then, we have

$$\begin{aligned} \mathcal{N}_k &= \partial_t \left[\frac{i}{3} \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0}} \frac{ke^{it\Phi(\bar{k})}}{i\Phi(\bar{k})} v_{k_1} v_{k_2} v_{k_3} \right] \\ &\quad - \frac{i}{3} \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0}} \frac{ke^{it\Phi(\bar{k})}}{i\Phi(\bar{k})} (\partial_t v_{k_1} v_{k_2} v_{k_3} + v_{k_1} \partial_t v_{k_2} v_{k_3} + v_{k_1} v_{k_2} \partial_t v_{k_3}) \\ &=: \partial_t(\mathcal{N}_1)_k + (\mathcal{N}_2)_k. \end{aligned} \quad (1.10)$$

Note that this corresponds to integration by parts on $\int \mathcal{N}(t) dt$. With (1.10), we see that smooth solutions to (1.10) satisfy

$$v(t) = v_0 + \mathcal{N}_1(t) - \mathcal{N}_1(0) + \int_0^t \mathcal{N}_2(t') + \mathcal{R}(t') dt'. \quad (1.11)$$

In (1.10), both terms have $\Phi(\bar{k})(\neq 0)$ in the denominators, and this provides smoothing.

Now, suppose that we have $C_t H_x^s$ estimates on \mathcal{N}_1 , \mathcal{N}_2 , and \mathcal{R} in (1.11). For \mathcal{N}_2 and \mathcal{R} , we obtain smallness thanks to the time integration (for small t .) However, there is no small constant for \mathcal{N}_1 . Thus, we can not close the argument to obtain a contraction.

³In the following, we only deal with the focusing case, i.e. with the + sign in (1.1), since our analysis is local-in-time and thus the focusing/defocusing nature of the equation is irrelevant.

In order to fix this problem, we use the idea from Section 6 in [1]. The idea is to separate the low frequency part of the non-resonant part \mathcal{N} before differentiating by parts. Let $v^{(n)} = P_n v$, where P_n is the Dirichlet projection onto the frequencies $\{|k| \leq n\}$. Then, write $\mathcal{N} = \mathcal{N}^{(n)} + \mathcal{N}^{(-n)}$, where $\mathcal{N}^{(n)}$ is given by

$$\mathcal{N}^{(n)} = \frac{i}{3} \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0}} k e^{it\Phi(\bar{k})} v_{k_1}^{(n)} v_{k_2}^{(n)} v_{k_3}^{(n)} \quad (1.12)$$

and $\mathcal{N}^{(-n)} = \mathcal{N} - \mathcal{N}^{(n)}$. Differentiating $\mathcal{N}^{(-n)}$ by parts, we obtain

$$\mathcal{N}^{(-n)} = \partial_t (\mathcal{N}_1^{(-n)}) + \mathcal{N}_2^{(-n)}, \quad (1.13)$$

where $\mathcal{N}_1^{(-n)}$ and $\mathcal{N}_2^{(-n)}$ are as in (1.10) with an extra condition

$$k^* := \max(|k_1|, |k_2|, |k_3|) > n. \quad (1.14)$$

Hence, smooth solutions to (1.10) satisfy

$$v(t) = v_0 + \mathcal{N}_1^{(-n)}(t) - \mathcal{N}_1^{(-n)}(0) + \int_0^t \mathcal{R}(t') + \mathcal{N}^{(n)}(t') + \mathcal{N}_2^{(-n)}(t') dt'. \quad (1.15)$$

It turns out that (1.14) provides a small constant $n^{-\alpha}$ for some $\alpha > 0$ in estimating $\mathcal{N}_1^{(-n)}$ (see Lemma 2.4 below), and we can close the argument for $s > \frac{1}{2}$.

Unfortunately, this turns out not to be sufficient when $s = \frac{1}{2}$. In particular, the estimate on $\mathcal{N}_2^{(-n)}$ fails when $s = \frac{1}{2}$. (See Lemma 2.3 below.) Thus, we need to proceed one step further. We hoped to take a differentiation by parts once more. A direct differentiation by parts, however, does not work because the corresponding resonance function $\Phi(\bar{k}) + \Phi(\bar{j})$ does not have a good factorization. (See (2.11) below.) In this case, we restrict $\mathcal{N}_2^{(-n)}$ into a part and then perform differentiation by parts once more, but in a slightly more complicated manner. See (4.10) and (4.17). Namely, in (1.10) and [1], we perform differentiation by parts to simply move the time derivative from a complex exponential to a product of v_{k_j} . However, in (4.17), we need to perform integration by parts⁴ to move the time derivative \mathcal{N}_1 (= product of $e^{it\Phi(\bar{j})}$ and $v_{j_1} v_{j_2} v_{j_3}$) to $e^{it\Phi(\bar{k})}$ and $v_{k_2} v_{k_3}$, which leads to further quinti- and septi-linear estimates. See Section 4 for details.

Lastly, we point out that the restriction $s \geq \frac{1}{2}$ on the regularity is due to the resonant term \mathcal{R} (see [2, p.228] and Lemma 2.1.) As pointed out in [12], if we define v by

$$v_k(t) = e^{ik^3 t + ik \int_0^t |u_k(t')|^2 t'} u_k(t), \quad (1.16)$$

then this would formally eliminate the resonant term. However, it is difficult to make sense of this transformation for nonsmooth functions. Instead, following [12, 11], one may try to use

$$v_k(t) = e^{ik^3 t + ik|u_k(0)|^2 t} u_k(t) \quad (1.17)$$

as the first order approximation to (1.16) in order to weaken the nonlinear effect of the resonant term \mathcal{R} . For further improvement, one may consider the second order approximation: $v_k(t) = e^{ik^3 t + ik(|u_k(0)|^2 t + \frac{1}{2} \partial_t |u_k(0)|^2 t^2)} u_k(t)$ or higher order approximations.

This paper is organized as follows. In Section 2, we prove a priori estimates needed for $s > \frac{1}{2}$. Then, we present the proof of Theorem 1.1 for $s > \frac{1}{2}$ in Section 3. In Section 4, we present the argument for the endpoint case $s = \frac{1}{2}$.

⁴Indeed, we keep (4.17) in the form of integration by parts to emphasize this point.

Acknowledgments: The authors would like to thank the anonymous referee for helpful suggestions.

2. NONLINEAR ESTIMATES FOR $s > \frac{1}{2}$

In this section, we present nonlinear estimates controlling the terms in (1.15). Without loss of generality, we assume that v_k is nonnegative in the following.

Lemma 2.1. *Let \mathcal{R} be as in (1.7). Then, for any $s \in \mathbb{R}$, we have*

$$\|\mathcal{R}(v)\|_{H^s} \lesssim \|v\|_{H^{\frac{1}{2}}}^2 \|v\|_{H^s}. \quad (2.1)$$

Also, for $s \geq \frac{1}{2}$, we have

$$\|\mathcal{R}(v) - \mathcal{R}(w)\|_{H^s} \lesssim (\|v\|_{H^s} + \|w\|_{H^s})^2 \|v - w\|_{H^s}. \quad (2.2)$$

Proof. We only prove (2.1) since (2.2) follows in a similar manner. Clearly, we have

$$\|\mathcal{R}(v)\|_{H^s} = \left(\sum_k |k|^{2+2s} v_k^6 \right)^{\frac{1}{2}} \leq \left\| |k|^{\frac{1}{2}} v_k \right\|_{l_k^\infty}^2 \|v\|_{H^s},$$

which is bounded by RHS of (2.1). \square

In the following, fix $n \in \mathbb{N}$.

Lemma 2.2. *Let $\mathcal{N}^{(n)}$ be as in (1.12). Then, for $s \geq \frac{1}{2}$, we have*

$$\|\mathcal{N}^{(n)}(v)\|_{H^s} \lesssim n \ln n \|v\|_{H^s}^3 \quad (2.3)$$

$$\|\mathcal{N}^{(n)}(v) - \mathcal{N}^{(n)}(w)\|_{H^s} \lesssim n \ln n (\|v\|_{H^s} + \|w\|_{H^s})^2 \|v - w\|_{H^s}. \quad (2.4)$$

Proof. We only prove (2.3) since (2.4) follows in a similar manner. Without loss of generality, assume $|k_1| \gtrsim |k|$. Then, by $|k| \lesssim n$ and Young's inequality, we have

$$\begin{aligned} \|\mathcal{N}^{(n)}(v)\|_{H^s} &\leq \left(\sum_k |k|^{2+2s} \left(\sum_{k_1+k_2+k_3=k} v_{k_1}^{(n)} v_{k_2}^{(n)} v_{k_3}^{(n)} \right)^2 \right)^{\frac{1}{2}} \\ &\lesssim n \left(\sum_k \left(\sum_{k_1+k_2+k_3=k} |k_1|^s v_{k_1}^{(n)} v_{k_2}^{(n)} v_{k_3}^{(n)} \right)^2 \right)^{\frac{1}{2}} \\ &\leq n \|v^{(n)}\|_{H^s} \|v_{k_2}^{(n)}\|_{l_{k_2}^1} \|v_{k_3}^{(n)}\|_{l_{k_3}^1} \lesssim n \ln n \|v^{(n)}\|_{H^s} \|v^{(n)}\|_{H^{\frac{1}{2}}}^2 \end{aligned}$$

where we used Cauchy-Schwarz inequality and $|k_j| \leq n$ in the last step. \square

Recall the following [2, (8.21), (8.22)]: Suppose $\Phi(\bar{k}) \neq 0$ when $k = k_1 + k_2 + k_3$. Then, we have the following two possibilities:

(a) With $k^* = \max(|k_1|, |k_2|, |k_3|)$,

$$\begin{aligned} |\Phi(k)| &\geq \max(|k_1 + k_2||k_2 + k_3|, |k_2 + k_3||k_3 + k_1|, |k_3 + k_1||k_1 + k_2|) \\ &\gtrsim (k^*)^2, \end{aligned} \quad (2.5)$$

In this case, we have $|\Phi(k)| \gtrsim (k^*)^2 \lambda$, where

$$\lambda = \lambda_k := \min(|k_1 + k_2|, |k_2 + k_3|, |k_3 + k_1|). \quad (2.6)$$

(b) $|k_1| \sim |k_2| \sim |k_3| \sim k^*$ and

$$|\Phi(k)| \geq \max(|k_1 + k_2|, |k_2 + k_3|, |k_3 + k_1|) \gtrsim k^*. \quad (2.7)$$

In this case, we have $|\Phi(k)| \gtrsim k^* \Lambda$, where

$$\Lambda = \Lambda_k := \min(|k_1 + k_2||k_2 + k_3|, |k_2 + k_3||k_3 + k_1|, |k_3 + k_1||k_1 + k_2|). \quad (2.8)$$

Lemma 2.3. Let $\mathcal{N}_2^{(-n)}$ be as in (1.13). Then, for $s > \frac{1}{2}$, we have

$$\|\mathcal{N}_2^{(-n)}(v)\|_{H^s} \lesssim \|v\|_{H^s}^5 \quad (2.9)$$

$$\|\mathcal{N}_2^{(-n)}(v) - \mathcal{N}_2^{(-n)}(w)\|_{H^s} \lesssim (\|v\|_{H^s} + \|w\|_{H^s})^4 \|v - w\|_{H^s}. \quad (2.10)$$

The same estimates hold for \mathcal{N}_2 in (1.10).

Proof. We only prove (2.9) since (2.10) follows in a similar manner. From (1.7) and (1.10), we can separate $\mathcal{N}_2^{(-n)}$ into two parts:

$$\begin{aligned} (\mathcal{N}_2^{(-n)})_k &= - \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} \frac{kk_1 e^{it\Phi(\bar{k})}}{i\Phi(\bar{k})} |v_{k_1}|^2 v_{k_1} v_{k_2} v_{k_3} \\ &\quad + \frac{1}{3} \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} \frac{kk_1 e^{it(\Phi(\bar{k})+\Phi(\bar{j}))}}{i\Phi(\bar{k})} v_{j_1} v_{j_2} v_{j_3} v_{k_2} v_{k_3} =: (\mathcal{N}_{21}^{(-n)})_k + (\mathcal{N}_{22}^{(-n)})_k, \end{aligned} \quad (2.11)$$

where $\Phi(\bar{j}) := \Phi(k_1, j_1, j_2, j_3)$.

• **Part 1:** First, we estimate $\mathcal{N}_{21}^{(-n)}$. By duality, it suffices to prove

$$\sum_k \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} M_1 |u_{k_1}|^2 u_{k_1} u_{k_2} u_{k_3} z_k \lesssim \|u\|_{L^2}^5 \quad (2.12)$$

where $\|z\|_{L^2} = 1$ and M_1 is given by

$$M_1 = M_1(k_1, k_2, k_3) := \frac{|k|^{1+s} |k_1|}{|\Phi(\bar{k})| |k_1|^{3s} |k_2|^s |k_3|^s}.$$

By Cauchy-Schwarz inequality, we have

$$\begin{aligned} \text{LHS of (2.12)} &\leq \left(\sum_{k_1, k_2, k_3} |u_{k_1}|^6 |u_{k_2}|^2 |u_{k_3}|^2 \right)^{\frac{1}{2}} \left(\sum_k \sum_{k_1+k_2+k_3=k} M_1^2 z_k^2 \right)^{\frac{1}{2}} \\ &\leq \mathcal{M}_1 \|u\|_{L^2}^5, \end{aligned}$$

where $\mathcal{M}_1 = (\sum_k \sum_{k_1+k_2+k_3=k} M_1^2 z_k^2)^{\frac{1}{2}}$.

◦ Case 1.a: $\Phi(\bar{k})$ satisfies (2.5). In this case, we have

$$M_1^2 \sim \frac{1}{(k^*)^{2-2s} |k_1|^{6s-2} |k_2|^{2s} |k_3|^{2s} \lambda^2} \leq \frac{1}{(k_*)^{8s} \lambda^2}$$

where $k_* = \min(|k_1|, |k_2|, |k_3|)$. Thus, for $s > \frac{1}{8}$, we have $\mathcal{M}_1 \lesssim 1$ by summing over k_i ($\neq k_*$) in λ^{-2} , k_* for $(k_*)^{-8s}$, and then k for z_k^2 . Hence, (2.12) holds for $s > \frac{1}{8}$.

- Case 1.b: $\Phi(\bar{k})$ satisfies (2.7). In this case, we have

$$M_1^2 \sim \frac{1}{(k^*)^{8s-2}\Lambda^2} \leq \frac{1}{(k_*)^{8s-2}\Lambda^2}.$$

Thus, for $s \geq \frac{1}{4}$, we have $\mathcal{M}_1 \lesssim 1$ by summing over two frequencies in Λ^{-2} , and then k for z_k^2 . Hence, (2.12) holds for $s \geq \frac{1}{4}$.

Therefore, the estimates (2.9) and (2.10) for $\mathcal{N}_{21}^{(-n)}$ hold as long as $s \geq \frac{1}{4}$.

- Part 2:** Next, we estimate $\mathcal{N}_{22}^{(-n)}$. By duality, it suffices to prove

$$\sum_k \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} M_2 u_{j_1} u_{j_2} u_{j_3} u_{k_2} u_{k_3} z_k \lesssim \|u\|_{L^2}^5 \quad (2.13)$$

where $\|z\|_{L^2} = 1$ and M_2 is given by

$$M_2 = M_2(j_1, j_2, j_3, k_2, k_3) := \frac{|k|^{1+s}|k_1|}{|\Phi(\bar{k})||j_1|^s|j_2|^s|j_3|^s|k_2|^s|k_3|^s}.$$

By Cauchy-Schwarz inequality, we have

$$\begin{aligned} \text{LHS of (2.13)} &\leq \left(\sum_{j_1, j_2, j_3, k_2, k_3} |u_{j_1}|^2 |u_{j_2}|^2 |u_{j_3}|^2 |u_{k_2}|^2 |u_{k_3}|^2 \right)^{\frac{1}{2}} \left(\sum_k \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1}} M_2^2 z_k^2 \right)^{\frac{1}{2}} \\ &\leq \mathcal{M}_2 \|u\|_{L^2}^5, \end{aligned}$$

where $\mathcal{M}_2 = \left(\sum_k \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1}} M_2^2 z_k^2 \right)^{\frac{1}{2}}$. Without loss of generality, assume

$$|j_1| = \max(|j_1|, |j_2|, |j_3|) \gtrsim |k_1|. \quad (2.14)$$

Also, we assume

$$|k_1| = \max(|k_1|, |k_2|, |k_3|) \quad (2.15)$$

in the following, since this corresponds to the worst case.

- Case 2.a: $\Phi(\bar{k})$ satisfies (2.5). In this case, we have

$$M_2^2 \lesssim \frac{1}{|j_2|^{2s}|j_3|^{2s}|k_2|^{2s}|k_3|^{2s}\lambda^2}, \quad (2.16)$$

where λ is as in (2.6). (Note that λ is of no help if $\lambda = |k_2 + k_3|$ and $|k_1| \gg |j_2|, |j_3| \gg |k_2|, |k_3|$.) Thus, for $s > \frac{1}{2}$, we have $\mathcal{M}_2 \lesssim 1$ by summing over j_2, j_3, k_2, k_3 , and then k for z_k^2 . Hence, (2.13) holds for $s > \frac{1}{2}$.

- Case 2.b: $\Phi(\bar{k})$ satisfies (2.7). In this case, we have $k^* \sim |k_1| \sim |k_2| \sim |k_3| \gtrsim |k|$. Then, for $s \geq \frac{1}{2}$, we have

$$M_2^2 \lesssim \frac{|k_1|^2}{(k^*)^{2s}|j_1|^{2s}|j_2|^{2s}|j_3|^{2s}\Lambda^2} \lesssim \frac{1}{|j_2|^{2s}|j_3|^{2s}\Lambda^2}, \quad (2.17)$$

where Λ is as in (2.8). Thus, we have $\mathcal{M}_2 \lesssim 1$ by summing over k_2, k_3 for Λ^{-2} , j_2, j_3 , and then k for z_k^2 . Hence, (2.13) holds for $s > \frac{1}{2}$. \square

Lemma 2.4. Let $\mathcal{N}_1^{(-n)}$ be as in (1.13). Then, for $s > 0$, there exists $\alpha > 0$ such that

$$\|\mathcal{N}_1^{(-n)}(v)\|_{H^s} \lesssim n^{-\alpha} \|v\|_{H^s}^3 \quad (2.18)$$

$$\|\mathcal{N}_1^{(-n)}(v) - \mathcal{N}_1^{(-n)}(w)\|_{H^s} \lesssim n^{-\alpha} (\|v\|_{H^s} + \|w\|_{H^s})^2 \|v - w\|_{H^s}. \quad (2.19)$$

Proof. We only prove (2.18) since (2.19) follows in a similar manner. Recall that

$$\mathcal{N}_1^{(-n)} = \frac{i}{3} \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} \frac{ke^{it\Phi(\bar{k})}}{i\Phi(\bar{k})} v_{k_1} v_{k_2} v_{k_3}.$$

By duality, it suffices to prove

$$\sum_k \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} M_3 u_{k_1} u_{k_2} u_{k_3} z_k \lesssim n^{-\alpha} \|u\|_{L^2}^3 \quad (2.20)$$

where $\|z\|_{L^2} = 1$ and M_3 is given by

$$M_3 = M_3(k_1, k_2, k_3) := \frac{|k|^{1+s}}{|\Phi(\bar{k})| |k_1|^s |k_2|^s |k_3|^s}.$$

As before, by Cauchy-Schwarz inequality, we have

$$\text{LHS of (2.20)} \leq \mathcal{M}_3 \|u\|_{L^2}^3,$$

where $\mathcal{M}_3 = (\sum_k \sum_{k_1+k_2+k_3=k} M_3^2 z_k^2)^{\frac{1}{2}}$. Without loss of generality, assume $k^* = |k_1|$.

o Case 1: $\Phi(\bar{k})$ satisfies (2.5). In this case, we have

$$M_3^2 \lesssim \frac{1}{(k^*)^2 |k_2|^{2s} |k_3|^{2s} \lambda^2} \leq n^{-1} \frac{1}{|k_2|^{1+2s} \lambda^2},$$

where λ is as in (2.6). Thus, for $s > 0$, we have $\mathcal{M}_3 \lesssim n^{-1}$ by summing over $k_j (\neq k_2)$ appearing in λ^{-2} , k_2 , and then k for z_k^2 . Hence, (2.20) holds for $s > 0$.

o Case 2: $\Phi(\bar{k})$ satisfies (2.7). For Λ as in (2.8), we have $\max(|k_2|, |k_3|, \Lambda) \gtrsim |k_1| > n$, since we have $|k_2| \sim |k_1|$ or $|k_3| \sim |k_1|$ if $\Lambda \ll |k_1|$. Then, we have

$$M_3^2 \lesssim \frac{1}{|k_2|^{2s} |k_3|^{2s} \Lambda^2} \lesssim \max(n^{-2s}, n^{-1+\varepsilon}) \frac{1}{\Lambda^{1+\varepsilon}},$$

Thus, for $s > 0$, we have $\mathcal{M}_3 \lesssim n^{-\alpha}$ by summing over two frequencies for $\Lambda^{-1-\varepsilon}$, and then k for z_k^2 . Hence, (2.20) holds for $s > 0$. \square

3. UNCONDITIONAL LOCAL WELL-POSEDNESS FOR $s > \frac{1}{2}$

In this section, we put together all the lemmata in the previous section and prove unconditional local well-posedness of mKdV (with prescribed L^2 -norm) in $H^s(\mathbb{T})$, $s > \frac{1}{2}$. Some parts of the argument below are standard. However, we include them for completeness.

For $n \in \mathbb{N}$, define $F_n(v, v_0)$ by $F_n(v, v_0) = F_n^{(1)}(v, v_0) + F_n^{(2)}(v, v_0)$, where $F_n^{(1)}$ and $F_n^{(2)}$ are given by

$$F_n^{(1)} = \mathcal{N}_1^{(-n)}(v)(t) - \mathcal{N}_1^{(-n)}(v_0)$$

$$F_n^{(2)} = \int_0^t \mathcal{R}(v)(t') + \mathcal{N}^{(n)}(v)(t') + \mathcal{N}_2^{(-n)}(v)(t') dt',$$

Then, if v is a solution to (1.7), then we have

$$v(t) = v_0 + F_n(v, v_0)(t). \quad (3.1)$$

Given an initial condition $v_0 \in H^s$, $s > \frac{1}{2}$, take a sequence $\{v_0^{[m]}\}_{m \in \mathbb{N}}$ of smooth functions such that $v_0^{[m]} \rightarrow v_0$ in H^s . Let $R = \|v_0\|_{H^s} + 1$. Then, without loss of generality, we can assume $\|v_0^{[m]}\|_{H^s} \leq R$.

Let $v^{[m]}$ denote the smooth global-in-time solution of mKdV with smooth initial condition $v_0^{[m]}$.⁵ Then, by Lemmata 2.1–2.4, we have

$$\begin{aligned} \|v^{[m]}\|_{C([0,T];H^s)} &\leq R + C \left\{ n^{-\alpha} (\|v^{[m]}\|_{C([0,T];H^s)} + \|v_0^{[m]}\|_{H^s})^2 \right. \\ &\quad \left. + n \ln n T (\|v^{[m]}\|_{C([0,T];H^s)}^2 + \|v^{[m]}\|_{C([0,T];H^s)}^4) \right\} \|v^{[m]}\|_{C([0,T];H^s)} \end{aligned} \quad (3.2)$$

First, choose n sufficiently large such that

$$Cn^{-\alpha}(3R)^2 < \frac{1}{4}. \quad (3.3)$$

Next, choose T sufficiently small such that

$$Cn \ln n T ((2R)^2 + (2R)^4) < \frac{1}{4}. \quad (3.4)$$

Then, from (3.2) with the continuity argument, we have

$$\|v^{[m]}\|_{C([0,T];H^s)} \leq 2R, \quad m \in \mathbb{N}.$$

Moreover, we have

$$\begin{aligned} &\|F_n(v^{[m_1]}, v_0^{[m_1]}) - F_n(v^{[m_2]}, v_0^{[m_2]})\|_{C([0,T];H^s)} \\ &\leq C_R \|v^{[m_1]} - v^{[m_2]}\|_{C([0,T];H^s)} + (1 + C_R) \|v_0^{[m_1]} - v_0^{[m_2]}\|_{H^s} \end{aligned} \quad (3.5)$$

where $C_R < \frac{1}{2}$ (by possibly taking larger n and smaller T .) Since $v^{[m_j]}$ is a (smooth) solution with initial condition $v_0^{[m_j]}$, it follows from (3.5) that

$$\|v^{[m_1]} - v^{[m_2]}\|_{C([0,T];H^s)} \leq C' \|v_0^{[m_1]} - v_0^{[m_2]}\|_{H^s} \quad (3.6)$$

for some $C' > 0$. Hence, $\{v^{[m]}\}$ converges in $C([0, T]; H^s)$.

Let v^∞ denote the limit. Then, we need to show that v^∞ satisfies (1.7) or

$$v(t) = v_0 + \int_0^t \mathcal{N}(v)(t') + \mathcal{R}(v)(t') dt' \quad (3.7)$$

as a space-time distribution. First, observe the following lemma. We present the proof at the end of this section.

Lemma 3.1. *Let \mathcal{N} and \mathcal{R} be as in (1.7). Then, we have, for any $\varepsilon > 0$,*

$$\|\mathcal{N} + \mathcal{R}\|_{H^{-\frac{1}{2}-\varepsilon}} \lesssim \|v\|_{H^{\frac{1}{2}}}^3. \quad (3.8)$$

In particular, if v satisfies (1.7), then we have $\|\partial_t v\|_{H^{-\frac{1}{2}-\varepsilon}} \lesssim \|v\|_{H^{\frac{1}{2}}}^3$.

⁵For smooth initial data, there exists a global smooth solution thanks to either the theory of complete integrability or the energy method. Instead of using smooth solutions, we could directly construct a solution by Galerkin approximation and compactness argument. Here, we use smooth solutions for conciseness of the presentation.

Given a test function ϕ , consider

$$\begin{aligned}
& \int_0^T \int_{\mathbb{T}} \left\{ v^\infty(t) - v_0 - \int_0^t \tilde{\mathcal{N}}(v^\infty)(t') dt' \right. \\
& \quad \left. - \left[v^{[m]}(t) - v_0^{[m]} - \int_0^t \tilde{\mathcal{N}}(v^{[m]})(t') dt' \right] \right\} \phi(x, t) dx dt \\
&= \int_0^T \int_{\mathbb{T}} (v^\infty(t) - v^{[m]}(t)) \phi dx dt - \int_0^T \int_{\mathbb{T}} (v_0 - v_0^{[m]}) \phi dx dt \\
& \quad + \int_0^T \int_{\mathbb{T}} \int_0^t \left[\tilde{\mathcal{N}}(v^\infty)(t') - \tilde{\mathcal{N}}(v^{[m]})(t') \right] dt' \phi(x, t) dx dt \\
&=: I_1 - I_2 + I_3
\end{aligned} \tag{3.9}$$

where $\tilde{\mathcal{N}} = \mathcal{N} + \mathcal{R}$. By convergence of $v^{[m]} \rightarrow v^\infty$ in $C([0, T]; H^s)$ and $v_0^{[m]} \rightarrow v_0$ in H^s , we have $I_1, I_2 \rightarrow 0$ as $m \rightarrow \infty$. By Lemma 3.1, we have

$$|I_3| \lesssim T(\|v^\infty\|_{C([0, T]; H^s)}^2 + \|v^{[m]}\|_{C([0, T]; H^s)}^2) \|v^\infty - v^{[m]}\|_{C([0, T]; H^s)} \|\phi\|_{L_t^1 H_x^{\frac{1}{2}+\varepsilon}} \rightarrow 0$$

as $m \rightarrow \infty$. Therefore, v^∞ is a solution to (3.7). It follows from (3.3) and (3.4) that the time of existence T satisfies $T \gtrsim \|v_0\|_{H^s}^{-\beta}$ for some $\beta > 0$. Also, the Lipschitz dependence on initial data follows from (3.6).

Let T be given. Suppose that both v and \tilde{v} are solutions in $C([0, T]; H^s)$ to (1.7) with the same initial condition $v_0 \in H^s(\mathbb{T})$, $s > \frac{1}{2}$. First, assume that $\|v\|_{C([0, T]; H^s)}, \|\tilde{v}\|_{C([0, T]; H^s)} \leq 2R$ where $R = \|v_0\|_{H^s} + 1$. Choose n and τ satisfying (3.3) and (3.4) (in place of T .) Then, from (3.5), we have

$$\|v - \tilde{v}\|_{C([0, \tau]; H^s)} = \|F_n(v, v_0) - F_n(\tilde{v}, v_0)\|_{C([0, \tau]; H^s)} \leq \frac{1}{2} \|v - \tilde{v}\|_{C([0, \tau]; H^s)}.$$

Hence, $v = \tilde{v}$ in $C([0, \tau]; H^s)$. By iterating the argument, we obtain $v = \tilde{v}$ in $C([0, T]; H^s)$. Now, suppose that $\tilde{R} := \frac{1}{2} \max(\|v\|_{C([0, T]; H^s)}, \|\tilde{v}\|_{C([0, T]; H^s)}) > R$. Then, use \tilde{R} (in place of R) to determine n and τ (in place of T) in (3.3) and (3.4). The rest follows as before. This proves the unconditional uniqueness (with prescribed L^2 -norm.)

We conclude this section by presenting the proof of Lemma 3.1.

Proof of Lemma 3.1. The contribution from \mathcal{R} is bounded by Lemma 2.1. Without loss of generality, assume $|k_1| = \max(|k_1|, |k_2|, |k_3|) \gtrsim |k|$. Then, by Young's inequality, we have

$$\begin{aligned}
\|\mathcal{N}\|_{H^{-\frac{1}{2}-\varepsilon}} &\lesssim \left(\sum_k \left(\sum_{k_1+k_2+k_3=k} |k_1|^{\frac{1}{2}} v_{k_1} |k_2|^{-\frac{1}{2}\varepsilon} v_{k_2} |k_2|^{-\frac{1}{2}\varepsilon} v_{k_3} \right)^2 \right)^{\frac{1}{2}} \\
&\leq \|v\|_{H^{\frac{1}{2}}} \| |k|^{-\frac{1}{2}\varepsilon} v_k \|_{l_k^1}^2 \lesssim \|v\|_{H^{\frac{1}{2}}}^3.
\end{aligned}$$

where we used Cauchy-Schwarz inequality in the last step. \square

4. ENDPOINT CASE: $s = \frac{1}{2}$

The previous argument fails at the endpoint regularity $s = \frac{1}{2}$ precisely because $\mathcal{N}_2^{(-n)}$ does not satisfy the required estimate when $s = \frac{1}{2}$. See Lemma 2.3. However, when $s = \frac{1}{2}$, Lemma 2.3 still holds for $\mathcal{N}_{21}^{(-n)}$ defined in (2.11). Moreover, if any of the following

conditions holds, then Lemma 2.3 holds for $\mathcal{N}_{22}^{(-n)}$ defined in (2.11), even when $s = \frac{1}{2}$:

- (a) $\max(|k_2|, |k_3|) \gtrsim \min(|k|^{\frac{1}{100}}, |k_1|^{\frac{1}{100}})$, if $\Phi(\bar{k})$ satisfies (2.5),
 - (b) $\max(|k_2|, |k_3|) \gtrsim \min(|j_2|^{\frac{1}{100}}, |j_3|^{\frac{1}{100}})$, if $\Phi(\bar{k})$ satisfies (2.5),
 - (c) $|j_1| \gtrsim \min(|k|^{1+\frac{1}{100}}, |k_1|^{1+\frac{1}{100}})$,
 - (d) $\Lambda_j \lesssim \max(|k_2|^{\frac{1}{100}}, |k_3|^{\frac{1}{100}})$,
- if $\Phi(\bar{k})$ satisfies (2.5) and $\Phi(\bar{j})$ satisfies (2.7),

$$(e) \quad \Lambda_k \gtrsim \max(|j_2|^{\frac{1}{100}}, |j_3|^{\frac{1}{100}}), \text{ if } \Phi(\bar{k}) \text{ satisfies (2.7)},$$

where Λ_j is as in (2.8) with $\{j_i\}$ in place of $\{k_i\}$.

By symmetry, assume $|k_2| \geq |k_3|$ and $|j_2| \geq |j_3|$. First, suppose that (a) holds. Then, in (2.16), we have

$$M_2^2 \lesssim \frac{1}{|j_1|^{2\varepsilon} |j_2|^{2s} |j_3|^{2s} |k_2|^{2s-200\varepsilon} |k_3|^{2s} \lambda^2} \lesssim \frac{1}{|j_2|^{2s+\varepsilon} |j_3|^{2s+\varepsilon} |k_2|^{2s-201\varepsilon} |k_3|^{2s+\varepsilon} \lambda^2}. \quad (4.1)$$

Thus, for $s = \frac{1}{2}$, we have $\mathcal{M}_2 \lesssim 1$ by summing over $k_i (\neq k_3)$ for λ^{-2} , k_3 , j_2, j_3 , and then k for z_k^2 . Hence, (2.13) holds for $s = \frac{1}{2}$. Secondly, suppose that (b) holds. By $|k_2|^{200\varepsilon} \gtrsim |j_2|^\varepsilon |j_3|^\varepsilon$, we have $M_2^2 \lesssim \text{RHS of (4.1)}$ as before. Next, suppose that (c) holds. Then, we have a small additional power of $|j_1|$ in the denominators of (2.16) and (2.17). Hence, (2.13) holds for $s = \frac{1}{2}$.

Now, suppose (d) holds. This implies that either

- (d.1) $|j_2 + j_3| \lesssim \max(|k_2|^{\frac{1}{100}}, |k_3|^{\frac{1}{100}})$, or
- (d.2) $|j_1 + j_2| |j_1 + j_3| \lesssim \max(|k_2|^{\frac{1}{100}}, |k_3|^{\frac{1}{100}})$

By symmetry, assume $|k_2| \geq |k_3|$. If (d.1) holds, then for fixed j_3 , there are at most $O(|k_2|^{\frac{1}{100}})$ possible choices for j_2 . Then, by going back to Case 2.a in Lemma 2.3, we have

$$M_2^2 \lesssim \frac{1}{|j_2|^{2s} |j_3|^{2s} |k_2|^{2s} |k_3|^{2s} \lambda^2} \lesssim \frac{1}{|j_3|^{4s} |k_2|^{\frac{1}{100}} |k_3|^{4s-\frac{1}{100}} \lambda^2}.$$

Thus, for $s > \frac{101}{400}$, we have $\mathcal{M}_2 \lesssim 1$ by summing over $k_i (\neq k_3)$ for λ^{-2} , k_3 , then, j_2 and j_3 , and finally k for z_k^2 . Hence, (2.13) holds for $s = \frac{1}{2}$.

If (d.2) holds, then for fixed j_1 , there are at most $O(|k_2|^{\frac{1}{100}})$ possible choices for j_3 . Since $\Phi(\bar{j})$ satisfies (2.7), we have $|j_1| \sim |j_2|$. Then, by going back to Case 2.a in Lemma 2.3, we have

$$M_2^2 \lesssim \frac{1}{|j_1|^{2s} |j_3|^{2s} |k_2|^{2s} |k_3|^{2s} \lambda^2} \lesssim \frac{1}{|j_1|^{4s} |k_2|^{\frac{1}{100}} |k_3|^{4s-\frac{1}{100}} \lambda^2}.$$

Once again, for $s > \frac{101}{400}$, we have $\mathcal{M}_2 \lesssim 1$ by summing over $k_i (\neq k_3)$ for λ^{-2} , k_3 , then, j_3 and j_1 , and finally k for z_k^2 . Hence, (2.13) holds for $s = \frac{1}{2}$.

Finally, suppose (e) holds. Then, in (2.17), we have

$$M_2^2 \lesssim \frac{1}{|j_2|^{2s+\varepsilon} |j_3|^{2s+\varepsilon} \Lambda_k^{2-200\varepsilon}}.$$

Thus, for $s = \frac{1}{2}$, we have $\mathcal{M}_2 \lesssim 1$ by summing over two frequencies for Λ_k , then, j_2 and j_3 , and finally k for z_k^2 . Hence, (2.13) holds for $s = \frac{1}{2}$.

Hence, letting $\mathcal{N}_{221}^{(-n)}$ be the restriction of $\mathcal{N}_{22}^{(-n)}$ such that at least one of the above conditions (a), (b), (c), (d), or (e) holds, we have the following estimates.

Lemma 4.1. *There exists $s_0 < \frac{1}{2}$ such that the following estimates hold for $s > s_0$:*

$$\|\mathcal{N}_{221}^{(-n)}(v)\|_{H^s} \lesssim \|v\|_{H^s}^5 \quad (4.2)$$

$$\|\mathcal{N}_{221}^{(-n)}(v) - \mathcal{N}_{221}^{(-n)}(w)\|_{H^s} \lesssim (\|v\|_{H^s} + \|w\|_{H^s})^4 \|v - w\|_{H^s}. \quad (4.3)$$

Now, letting $\mathcal{N}_{222}^{(-n)} := \mathcal{N}_{22}^{(-n)} - \mathcal{N}_{221}^{(-n)}$, we have

$$\mathcal{N}_{22}^{(-n)} = \mathcal{N}_{221}^{(-n)} + \mathcal{N}_{222}^{(-n)}. \quad (4.4)$$

In the following, we concentrate on estimating the contribution from $\mathcal{N}_{222}^{(-n)}$. Note that $\mathcal{N}_{222}^{(-n)}$ is the restriction of $\mathcal{N}_{22}^{(-n)}$ such that *all* of the conditions below hold:

$$(a') \quad \max(|k_2|, |k_3|) \ll \min(|k|^{\frac{1}{100}}, |k_1|^{\frac{1}{100}}), \quad (4.5)$$

$$(b') \quad \max(|k_2|, |k_3|) \ll \min(|j_2|^{\frac{1}{100}}, |j_3|^{\frac{1}{100}}), \quad (4.6)$$

$$(c') \quad |j_1| \ll \min(|k|^{1+\frac{1}{100}}, |k_1|^{1+\frac{1}{100}}), \quad (4.7)$$

$$(d') \quad \Lambda_j \gg \max(|k_2|^{\frac{1}{100}}, |k_3|^{\frac{1}{100}}), \quad (4.8)$$

$$(e') \quad \Lambda_k \ll \max(|j_2|^{\frac{1}{100}}, |j_3|^{\frac{1}{100}}), \quad (4.9)$$

where (a') and (b') hold when $\Phi(\bar{k})$ satisfies (2.5), (d') holds when $\Phi(\bar{k})$ satisfies (2.5) and $\Phi(\bar{j})$ satisfies (2.7), and (e') holds when $\Phi(\bar{k})$ satisfies (2.7). (Recall that we also assume (2.14)–(2.15).) Henceforth, we assume that the frequencies are restricted such that the conditions (a')–(e') hold. By (1.7) and (1.10), we have

$$\begin{aligned} (\mathcal{N}_{222}^{(-n)})_k &= \frac{1}{3} \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} \frac{kk_1 e^{it(\Phi(\bar{k}) + \Phi(\bar{j}))}}{i\Phi(\bar{k})} v_{j_1} v_{j_2} v_{j_3} v_{k_2} v_{k_3} \\ &= -i \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} \frac{ke^{it(\Phi(\bar{k}))}}{i\Phi(\bar{k})} \partial_t(\mathcal{N}_1)_{k_1} v_{k_2} v_{k_3} + i \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} \frac{ke^{it(\Phi(\bar{k}))}}{i\Phi(\bar{k})} (\mathcal{N}_2)_{k_1} v_{k_2} v_{k_3} \\ &=: (\mathcal{N}_3)_k + (\mathcal{N}_4)_k. \end{aligned} \quad (4.10)$$

The following lemma shows that \mathcal{N}_4 can be controlled in $H^{\frac{1}{2}}$.

Lemma 4.2. *The following estimates hold:*

$$\|\mathcal{N}_4(v)\|_{H^{\frac{1}{2}}} \lesssim \|v\|_{H^{\frac{1}{2}}}^7 \quad (4.11)$$

$$\|\mathcal{N}_4(v) - \mathcal{N}_4(w)\|_{H^{\frac{1}{2}}} \lesssim (\|v\|_{H^{\frac{1}{2}}} + \|w\|_{H^{\frac{1}{2}}})^6 \|v - w\|_{H^{\frac{1}{2}}}. \quad (4.12)$$

Before proving this lemma, let us present the following corollary to Lemma 2.3.

Corollary 4.3. *For $s < \frac{1}{2}$, we have*

$$\|\mathcal{N}_2(v)\|_{H^s} \lesssim \|v\|_{H^{\frac{1}{2}}}^5 \quad (4.13)$$

$$\|\mathcal{N}_2(v) - \mathcal{N}_2(w)\|_{H^s} \lesssim (\|v\|_{H^{\frac{1}{2}}} + \|w\|_{H^{\frac{1}{2}}})^4 \|v - w\|_{H^{\frac{1}{2}}}. \quad (4.14)$$

We omit the proof of this corollary, since it follows from a slight modification of the proof of Lemma 2.3.

Proof of Lemma 4.2. We only prove (4.11) since (4.12) follows in a similar manner. In view of Corollary 4.3, it suffices to prove

$$\|\mathcal{N}_4\|_{H^{\frac{1}{2}}} \lesssim \|\mathcal{N}_2\|_{H^{-\frac{1}{2}}} \|v\|_{H^{\frac{1}{2}}}^2. \quad (4.15)$$

By duality, it suffices to prove

$$\sum_k \sum_{\substack{k_1+k_2+k_3=k \\ \Phi(\bar{k}) \neq 0 \\ k^* > n}} M_4 u_{k_1} u_{k_2} u_{k_3} z_k \lesssim \|u\|_{L^2}^3 \quad (4.16)$$

where $\|z\|_{L^2} = 1$ and M_4 is given by

$$M_4 = M_4(k_1, k_2, k_3) := \frac{|k|^{\frac{3}{2}} |k_1|^{\frac{1}{2}}}{|\Phi(\bar{k})| |k_2|^{\frac{1}{2}} |k_3|^{\frac{1}{2}}}.$$

As before, by Cauchy-Schwarz inequality, we have LHS of (4.16) $\leq \mathcal{M}_4 \|u\|_{L^2}^3$, where $\mathcal{M}_4 = (\sum_k \sum_{k_1+k_2+k_3=k} M_4^2 z_k^2)^{\frac{1}{2}}$.

- o Case 1: $\Phi(\bar{k})$ satisfies (2.5). In this case, we have $M_4^2 \lesssim |k_2|^{-1} |k_3|^{-1} \lambda^{-2} \leq |k_3|^{-2} \lambda^{-2}$, where λ is as in (2.6). Thus, we have $\mathcal{M}_4 \lesssim 1$ by summing over k_i ($\neq k_3$) appearing in λ^{-2} , k_3 , and then k for z_k^2 . Hence, (4.16) holds.
- o Case 2: $\Phi(\bar{k})$ satisfies (2.7). By $|k_2| \sim |k_1| \sim k^*$, we have $M_4^2 \lesssim \Lambda^{-2}$ where Λ is as in (2.8). Thus, we have $\mathcal{M}_4 \lesssim 1$ by summing over two frequencies for Λ^{-2} , and then k for z_k^2 . Hence, (4.16) holds. \square

Now, we need to estimate the contribution from \mathcal{N}_3 . In doing so, we actually estimate $\int_0^t \mathcal{N}_3(t') dt'$. Integrating by parts, we have

$$\begin{aligned} \int_0^t \mathcal{N}_3(t') dt' &= -\frac{1}{3} \int_0^t \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} \frac{kk_1 e^{it'(\Phi(\bar{k})+\Phi(\bar{j}))}}{i\Phi(\bar{j})} (v_{j_1} v_{j_2} v_{j_3} v_{k_2} v_{k_3})(t') dt' \\ &\quad + \frac{1}{3} \int_0^t \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} \frac{kk_1 e^{it'(\Phi(\bar{k})+\Phi(\bar{j}))}}{\Phi(\bar{k})\Phi(\bar{j})} (v_{j_1} v_{j_2} v_{j_3})(t') \partial_t (v_{k_2} v_{k_3})(t') dt' \\ &\quad - \frac{1}{3} \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} \frac{kk_1 e^{it(\Phi(\bar{k})+\Phi(\bar{j}))}}{\Phi(\bar{k})\Phi(\bar{j})} (v_{j_1} v_{j_2} v_{j_3} v_{k_2} v_{k_3})(t') \Big|_0^t \\ &=: \int_0^t \mathcal{N}_5(t') dt' + \int_0^t \mathcal{N}_6(t') dt' + \mathcal{N}_7(v)(t) - \mathcal{N}_7(v)(0). \end{aligned} \quad (4.17)$$

First, note that \mathcal{N}_5 looks like $\mathcal{N}_2^{(-n)}$ in Lemma 2.3. However, it satisfies a better estimate thanks to the conditions (b')–(e').

Lemma 4.4. *The following estimates hold:*

$$\|\mathcal{N}_5(v)\|_{H^{\frac{1}{2}}} \lesssim \|v\|_{H^{\frac{1}{2}}}^5 \quad (4.18)$$

$$\|\mathcal{N}_5(v) - \mathcal{N}_5(w)\|_{H^{\frac{1}{2}}} \lesssim (\|v\|_{H^{\frac{1}{2}}} + \|w\|_{H^{\frac{1}{2}}})^4 \|v - w\|_{H^{\frac{1}{2}}}. \quad (4.19)$$

Proof. We only prove (4.18) since (4.19) follows in a similar manner. By duality, it suffices to prove

$$\sum_k \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} M_5 u_{j_1} u_{j_2} u_{j_3} u_{k_2} u_{k_3} z_k \lesssim \|u\|_{L^2}^5 \quad (4.20)$$

where $\|z\|_{L^2} = 1$ and M_5 is given by

$$M_5 = M_5(j_1, j_2, j_3, k_2, k_3) := \frac{|k|^{1+s}|k_1|}{|\Phi(\bar{j})||j_1|^s|j_2|^s|j_3|^s|k_2|^s|k_3|^s}$$

with $s = \frac{1}{2}$. As before, by Cauchy-Schwarz inequality, we have LHS of (4.20) $\leq \mathcal{M}_5 \|u\|_{L^2}^5$, where $\mathcal{M}_5 = (\sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1}} M_5^2 z_k^2)^{\frac{1}{2}}$.

By symmetry, assume $|k_2| \geq |k_3|$ and $|j_2| \geq |j_3|$.

◦ Case 1: Both $\Phi(\bar{k})$ and $\Phi(\bar{j})$ satisfy (2.5). Let λ_j be as in (2.6) with $\{j_i\}$ in place of $\{k_i\}$. Then, as in Case 2.a in Lemma 2.3, we have

$$M_5^2 \lesssim \frac{1}{|j_2|^{2s}|j_3|^{2s}|k_2|^{2s}|k_3|^{2s}\lambda_j^2} \ll \frac{1}{|j_3|^{4s-\frac{2\varepsilon}{100}}|k_2|^{2s+\varepsilon}|k_3|^{2s+\varepsilon}\lambda_j^2},$$

since $|j_2| \gg |k_2|^{100} \geq |k_3|^{100}$ by the condition (b'). Thus, for $s = \frac{1}{2}$, we have $\mathcal{M}_5 \lesssim 1$ by summing over k_2, k_3 , then $j_i (\neq j_3)$, for λ_j^{-2} , j_3 , and k for z_k^2 . Hence, (4.20) holds for $s = \frac{1}{2}$.

◦ Case 2: $\Phi(\bar{k})$ satisfies (2.7) and $\Phi(\bar{j})$ satisfies (2.5). By (e'), given k_3 , there are at most $O(|j_2|^{\frac{1}{50}})$ possible choices for k_2 . Then, by (c') with $|k_2| \sim |k_1|$, given k_3 , there are at most $O(|k_2|^{\frac{101}{5000}})$ possible choices for k_2 . As in Case 2.a in Lemma 2.3, we have

$$M_5^2 \lesssim \frac{1}{|j_2|^{2s}|j_3|^{2s}|k_2|^{2s}|k_3|^{2s}\lambda_j^2} \ll \frac{1}{|j_3|^{4s}|k_2|^{\frac{1}{40}}|k_3|^{4s-\frac{1}{40}}\lambda_j^2}.$$

Thus, for $s = \frac{1}{2}$, we have $\mathcal{M}_5 \lesssim 1$ by summing over k_2, k_3 , then $j_i (\neq j_3)$, for λ_j^{-2} , j_3 , and k for z_k^2 . Hence, (4.20) holds for $s = \frac{1}{2}$.

◦ Case 3: $\Phi(\bar{k})$ satisfies (2.5) and $\Phi(\bar{j})$ satisfies (2.7). In this case, we have $j^* \sim |j_1| \sim |j_2| \sim |j_3| \gtrsim |k_1|$. By (d') in (4.8), we have

$$M_5^2 \lesssim \frac{1}{|k_2|^{2s}|k_3|^{2s}\Lambda_j^2} \ll \frac{1}{|k_2|^{2s+\varepsilon}|k_3|^{2s+\varepsilon}\Lambda_j^{2-200\varepsilon}}$$

for $s \geq \frac{1}{2}$. Thus, we have $\mathcal{M}_2 \lesssim 1$ by summing over two frequencies for $\Lambda_j^{-2+200\varepsilon}$, k_2, k_3 , and then k for z_k^2 . Hence, (4.20) holds for $s = \frac{1}{2}$.

- Case 4: Both $\Phi(\bar{k})$ and $\Phi(\bar{j})$ satisfy (2.7). As in Case 2, given k_3 , there are at most $O(|k_2|^{\frac{1}{40}})$ possible choices for k_2 . Thus, we have

$$M_5^2 \lesssim \frac{1}{|k_2|^{2s} |k_3|^{2s} \Lambda_j^2} \ll \frac{1}{|k_2|^{\frac{1}{40}} |k_3|^{4s - \frac{1}{40}} \Lambda_j^{2-200\varepsilon}}$$

for $s \geq \frac{1}{2}$. Hence, (4.20) holds for $s = \frac{1}{2}$. \square

Lemma 4.5. *Suppose that v and w satisfy (1.7). Then, the following estimates hold:*

$$\|\mathcal{N}_6(v)\|_{H^{\frac{1}{2}}} \lesssim \|v\|_{H^{\frac{1}{2}}}^7 \quad (4.21)$$

$$\|\mathcal{N}_6(v) - \mathcal{N}_6(w)\|_{H^{\frac{1}{2}}} \lesssim (\|v\|_{H^{\frac{1}{2}}} + \|w\|_{H^{\frac{1}{2}}})^6 \|v - w\|_{H^{\frac{1}{2}}}. \quad (4.22)$$

Proof. Assume that the time derivative falls on v_{k_2} in (4.17). i.e. we have

$$\mathcal{N}_6(t) \sim \sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} \frac{kk_1 e^{it'(\Phi(\bar{k})+\Phi(\bar{j}))}}{\Phi(\bar{k})\Phi(\bar{j})} (v_{j_1} v_{j_2} v_{j_3}) \partial_t v_{k_2} v_{k_3}.$$

Moreover, assume $|j_2| \geq |j_3|$ and $|k_2| \geq |k_3|$. We only prove (4.21) since (4.22) follows in a similar manner. In view of Lemma 3.1, it suffices to prove

$$\|\mathcal{N}_6\|_{H^{\frac{1}{2}}} \lesssim \|\partial_t v\|_{H^{-\frac{1}{2}-\varepsilon}} \|v\|_{H^{\frac{1}{2}}}^4 \quad (4.23)$$

for some small $\varepsilon > 0$. By duality, it suffices to prove

$$\sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1 \\ \Phi(\bar{k}), \Phi(\bar{j}) \neq 0 \\ k^* > n}} M_6 u_{j_1} u_{j_2} u_{j_3} u_{k_2} u_{k_3} z_k \lesssim \|u\|_{L^2}^5 \quad (4.24)$$

where $\|z\|_{L^2} = 1$ and M_6 is given by

$$M_6 = M_6(j_1, j_2, j_3, k_2, k_3) := \frac{|k|^{\frac{3}{2}} |k_1| |k_2|^{\frac{1}{2}+\varepsilon}}{|\Phi(\bar{k})| |\Phi(\bar{j})| |j_1|^{\frac{1}{2}} |j_2|^{\frac{1}{2}} |j_3|^{\frac{1}{2}} |k_3|^{\frac{1}{2}}}.$$

As before, by Cauchy-Schwarz inequality, we have LHS of (4.24) $\leq \mathcal{M}_6 \|u\|_{L^2}^5$, where $\mathcal{M}_6 = (\sum_{\substack{k_1+k_2+k_3=k \\ j_1+j_2+j_3=k_1}} M_6^2 z_k^2)^{\frac{1}{2}}$. Let λ_k and λ_j (and Λ_k and Λ_j) be as in (2.6) (and as in (2.8)) for $\{k_i\}$ and $\{j_i\}$, respectively.

- Case 1: Both $\Phi(\bar{k})$ and $\Phi(\bar{j})$ satisfy (2.5). By (b') and $|j_1| \geq \max(|j_2|, |j_3|)$, we have

$$M_6^2 \lesssim \frac{1}{|j_1|^{4-\frac{1}{50}(\frac{1}{2}+\varepsilon)} |j_2| |j_3| |k_3| \lambda_k^2 \lambda_j^2} \leq \frac{1}{|j_1|^{1+\varepsilon} |j_2|^{1+\varepsilon} |j_3|^{1+\varepsilon} |k_3|^{1+\varepsilon} \lambda_k^2 \lambda_j^2}$$

for sufficiently small $\varepsilon > 0$. Thus, we have $\mathcal{M}_6 \lesssim 1$ by summing over k_3, j_1, j_2, j_3 , and then k for z_k^2 . Hence, (4.24) holds for $s = \frac{1}{2}$.

- Case 2: $\Phi(\bar{k})$ satisfies (2.5) and $\Phi(\bar{j})$ satisfies (2.7). In this case, we have

$$M_6^2 \lesssim \frac{1}{|j_1|^{2-\frac{1}{50}(\frac{1}{2}+\varepsilon)} |j_2| |j_3| |k_3| \lambda_k^2 \Lambda_j^2} \leq \frac{1}{|j_1|^{1+\varepsilon} |j_2|^{1+\varepsilon} |j_3|^{1+\varepsilon} |k_3|^{1+\varepsilon} \lambda_k^2 \Lambda_j^2}$$

for sufficiently small $\varepsilon > 0$. Thus, we have $\mathcal{M}_6 \lesssim 1$ by summing over k_3, j_1, j_2, j_3 , and then k for z_k^2 . Hence, (4.24) holds for $s = \frac{1}{2}$.

o Case 3: $\Phi(\bar{k})$ satisfies (2.7) and $\Phi(\bar{j})$ satisfies (2.5). In this case, we have

$$M_6^2 \lesssim \frac{1}{|j_1||j_2||j_3||k_3|^{1-2\varepsilon}\Lambda_k^2\lambda_j^2} \leq \frac{1}{|j_2|^{\frac{3}{2}}|j_3|^{\frac{3}{2}}|k_3|^{1-2\varepsilon}\Lambda_k^2\lambda_j^2}.$$

Thus, we have $\mathcal{M}_6 \lesssim 1$ by summing over two frequencies for Λ_k^{-2} , j_2, j_3 , and then k for z_k^2 . Hence, (4.24) holds for $s = \frac{1}{2}$.

o Case 4: Both $\Phi(\bar{k})$ and $\Phi(\bar{j})$ satisfy (2.7). In this case, we have $|k_1| \sim |k_2| \sim |k_3|$ and $|j_1| \sim |j_2| \sim |j_3|$. In this case, we have $M_6^2 \lesssim |j_1|^{-1+\varepsilon}\Lambda_k^{-2}\Lambda_j^{-2}$. Thus, we have $\mathcal{M}_6 \lesssim 1$ by summing over two frequencies for Λ_k^{-2} , two frequencies for Λ_j^{-2} , and then k for z_k^2 . Hence, (4.24) holds for $s = \frac{1}{2}$. \square

Before we present the estimate on \mathcal{N}_7 , recall that we have $|k_1| > n$. Thus, we have $|j_1| \gtrsim n$ since $|j_1| \gtrsim |k_1|$.

Lemma 4.6. *There exists $\alpha > 0$ such that*

$$\|\mathcal{N}_7(v)\|_{H^{\frac{1}{2}}} \lesssim n^{-\alpha} \|v\|_{H^{\frac{1}{2}}}^5 \quad (4.25)$$

$$\|\mathcal{N}_7(v) - \mathcal{N}_7(w)\|_{H^{\frac{1}{2}}} \lesssim n^{-\alpha} (\|v\|_{H^{\frac{1}{2}}} + \|w\|_{H^{\frac{1}{2}}})^4 \|v - w\|_{H^{\frac{1}{2}}}. \quad (4.26)$$

Proof. This lemma immediately follows from the proof of Lemma 4.5 once we note that there is an extra (small) power of $|j_1|$ in the denominator except for Case 3. In Case 3, we have an extra (small) power of $|k_3| \sim |k_1| > n$. \square

The remaining part of the argument is basically the same as in Section 3. For $n \in \mathbb{N}$, define $G_n(v, v_0)$ by $G_n(v, v_0) = G_n^{(1)}(v, v_0) + G_n^{(2)}(v, v_0)$, where $G_n^{(1)}$ and $G_n^{(2)}$ are given by

$$\begin{aligned} G_n^{(1)} &= \mathcal{N}_1^{(-n)}(v)(t) - \mathcal{N}_1^{(-n)}(v_0) + \mathcal{N}_7(v)(t) - \mathcal{N}_7(v_0) \\ G_n^{(2)} &= \int_0^t \mathcal{R}(v)(t') + \mathcal{N}^{(n)}(v)(t') + \mathcal{N}_{21}^{(-n)}(v)(t') \\ &\quad + \mathcal{N}_{221}^{(-n)}(v)(t') + \mathcal{N}_4(v)(t') + \mathcal{N}_5(v)(t') + \mathcal{N}_6(v)(t') dt'. \end{aligned}$$

From (1.15), (2.11), (4.4), (4.10), and (4.17), we have

$$v(t) = v_0 + G_n(v, v_0)(t).$$

if v is a solution to (1.7). It follows from Lemmata 2.1–2.4, 4.1, 4.2, and 4.4–4.6 that we obtain a priori bounds (3.2) and (3.5) for $s = \frac{1}{2}$ (with G_n in place of F_n and with a higher power in (3.2).) Then, the unconditional local well-posedness in $H^{\frac{1}{2}}$ with Lipschitz dependence on initial data follows as in Section 3.

REFERENCES

- [1] A. Babin, A. Ilyin, E. Titi, *On the regularization mechanism for the periodic Korteweg-de Vries equation*, Comm. Pure Appl. Math. 64 (2011), no. 5, 591–648.
- [2] J. Bourgain, *Fourier transform restriction phenomena for certain lattice subsets and applications to nonlinear evolution equations II*, Geom. Funct. Anal., 3 (1993), 209–262.
- [3] J. Bourgain, *Periodic Korteweg de Vries equation with measures as initial data*, Selecta Math. (N.S.) 3 (1997), no. 2, 115–159.
- [4] M. Christ, *Nonuniqueness of weak solutions of the nonlinear Schrödinger equation*, arXiv:math/0503366v1 [math.AP].
- [5] M. Christ, J. Colliander, T. Tao, *Asymptotics, frequency modulation, and low regularity ill-posedness for canonical defocusing equations*, Amer. J. Math. 125 (2003), no. 6, 1235–1293.

- [6] J. Colliander, M. Keel, G. Staffilani, H. Takaoka, T. Tao, *Sharp Global Well-Posedness for KdV and Modified KdV on \mathbb{R} and \mathbb{T}* , J. Amer. Math. Soc. 16 (2003), no. 3, 705–749.
- [7] G. Furioli, F. Planchon, E. Terraneo, *Unconditional well-posedness for semilinear Schrödinger and wave equations in H^s* , Harmonic analysis at Mount Holyoke (South Hadley, MA, 2001), 147–156, Contemp. Math., 320, Amer. Math. Soc., Providence, RI, 2003.
- [8] J. Ginibre, *An introduction to nonlinear Schrödinger equations*, Nonlinear waves (Sapporo, 1995), 85–133, GAKUTO Internat. Ser. Math. Sci. Appl., 10, Gakkōtoshō, Tokyo, 1997.
- [9] T. Kappeler, P. Topalov, *Global well-posedness of mKdV in $L^2(\mathbb{T}, \mathbb{R})$* , Comm. PDE 30 (2005), no. 1-3, 435–449.
- [10] T. Kato, *On nonlinear Schrödinger equations. II. H^s -solutions and unconditional well-posedness*, J. Anal. Math. 67 (1995), 281–306.
- [11] K. Nakanishi, H. Takaoka, Y. Tsutsumi, *Unique local existence of solution in low regularity space of the Cauchy problem for the mKdV equation with periodic boundary condition*, Séminaire: Équations aux Dérivées Partielles. 2007–2008, Exp. No. XVII, 7 pp., Sémin. Équ. Dériv. Partielles, École Polytech., Palaiseau, 2009.
- [12] H. Takaoka, Y. Tsutsumi, *Well-posedness of the Cauchy problem for the modified KdV equation with periodic boundary condition*, Int. Math. Res. Not. (2004), no. 56, 3009–3040.
- [13] T. Tao, *Global existence and uniqueness results for weak solutions of the focusing mass-critical nonlinear Schrödinger equation*, Anal. PDE 2 (2009), no. 1, 61–81.
- [14] Y. Zhou, *Uniqueness of weak solution of the KdV equation*, Int. Math. Res. Not. (1997), no. 6, 271–283.

SOONSIK KWON, DEPARTMENT OF MATHEMATICAL SCIENCES, KOREA ADVANCED INSTITUTE OF SCIENCE AND TECHNOLOGY, 335 GWAHANGNO, YUSEONG-GU, DAEJEON 305-701, REPUBLIC OF KOREA

E-mail address: soonsikk@kaist.edu

TADAHIRO OH, DEPARTMENT OF MATHEMATICS, UNIVERSITY OF TORONTO, 40 ST. GEORGE ST, TORONTO, ON M5S 2E4, CANADA

Current address: Department of Mathematics, Princeton University, Fine Hall, Washington Rd, Princeton, NJ 08544-1000, USA

E-mail address: hirooh@math.princeton.edu